

## Direct determination of mean-field from data on matter density

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In this work we consider the single particle Schrodinger equation and develop a method for determining the single particle potential  $V$  from a given single particle wave function  $\Psi(\vec{r})$  or matter density, assuming it is known for all  $\vec{r}$ . In particular, we consider the case of spherical symmetry. The results of this work are very important for investigation of the validity of the shell model and for the development of a modern EDF which provides enhanced predictive power for properties of nuclei and the equation of state (EOS) of nuclear matter (NM).

For the single particle Schrodinger equation,

$$-\frac{\hbar^2}{2m}\Delta\Psi + V\Psi = E\Psi, \quad (1)$$

we have that for a given single particle wave function  $\Psi(\vec{r})$ , known for all  $\vec{r}$ , the corresponding single particle potential  $V$  is uniquely determined [1] from

$$V(\vec{r}) = E + \frac{\hbar^2}{2m} \frac{\Delta\Psi(\vec{r})}{\Psi(\vec{r})}. \quad (2)$$

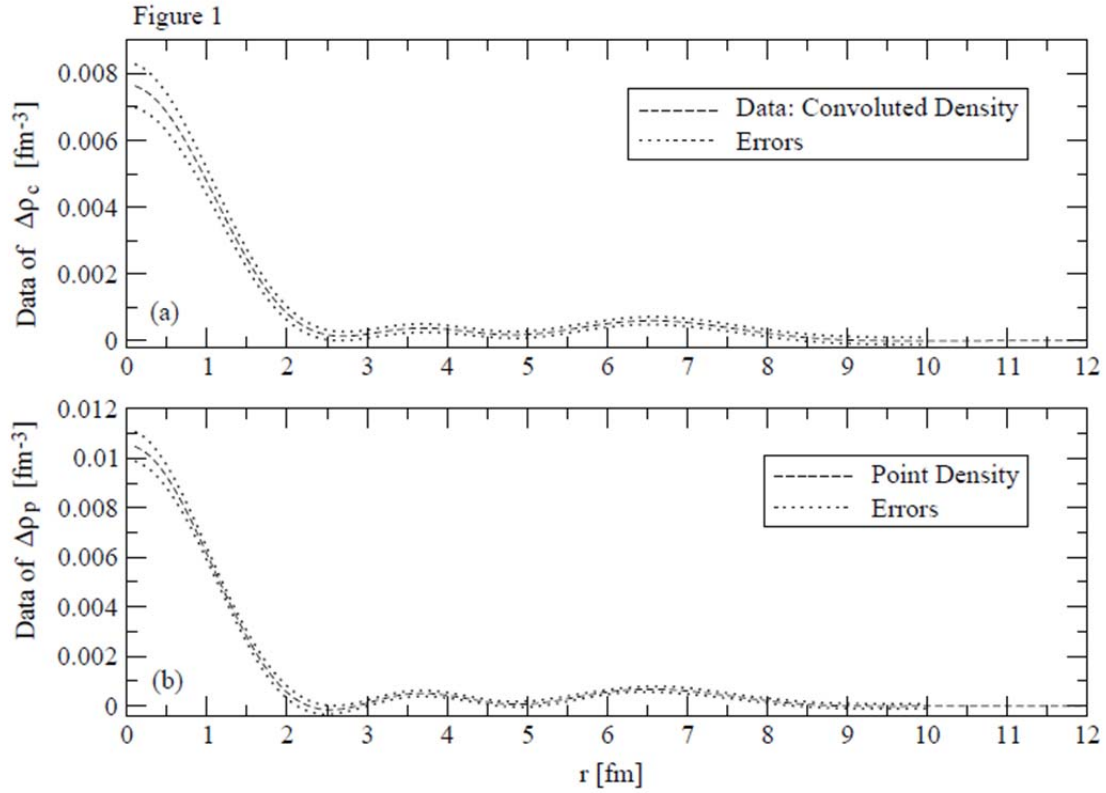
Note that for a nonsingular  $V$ ,  $\Delta\Psi(\vec{r}) = 0$  when  $\Psi(\vec{r}) = 0$ . In the spherical case we have for the centroid potential

$$V_{cen}(r) = E + \frac{\hbar^2}{2m} \frac{d^2 R_{nlj}}{dr^2} \frac{1}{R_{nlj}(r)} - \frac{\hbar^2}{2m} \frac{l(l+1)}{r} - \frac{1}{2}(1 - \tau_z)V_{coul}(r) - c_{ls} V_{s.o.}(r). \quad (3)$$

Here,  $V_{cen}(r)$ ,  $\vec{s} \cdot \vec{l} V_{s.o.}(r)$  and  $\frac{1}{2}(1 - \tau_z)V_{coul}(r)$ , are the central, spin-orbit and coulomb potentials, respectively, and  $\tau_z=1$  for a neutron and -1 for a proton.

We consider, in particular, the charge density difference,  $\Delta\rho_c(r)$ , between the isotones  $^{206}\text{Pb} - ^{205}\text{Tl}$ , associated with the proton  $3S_{1/2}$  single particle orbit, and determine the corresponding single particle potential. The experimental data for the charge densities,  $\rho_c(r)$ , of the isotones  $^{206}\text{Pb}$  and  $^{205}\text{Tl}$ , obtained from accurate elastic electron scattering experiments, are taken from Ref. [1]. In Fig. 1a we present the experimental data for the charge density difference,  $\Delta\rho_c(r)$ , between the isotones  $^{206}\text{Pb} - ^{205}\text{Tl}$ , shown by the dashed line. It is normalized to a total charge of one proton ( $Z=1$ ). The dotted lines indicate the experimental uncertainty. Note that the two nodes associated with the proton  $3S_{1/2}$  orbit are clearly seen in the figure. To extract the corresponding single particle potential we need the point proton distribution,  $\Delta\rho_p(r)$ . This is obtained by determining the point proton form factor,  $F_p(q)$ , and then extracting  $\Delta\rho_p(r)$ . The results are shown in Fig. 1b. Note that  $\Delta\rho_p(r)$  (dashed line) is slightly negative at the first node (at

$\sim 2.6$  fm) and above zero at the second node ( $r \sim 4.9$  fm). Using these results we determine the corresponding mean field associated with the  $3S_{1/2}$  orbit in  $^{206}\text{Pb}$  [2].



**FIG. 1.** Charge density (a) and point density (b) for the  $3S_{1/2}$  orbit in  $^{206}\text{Pb}$ .

[1] J.M. Cavedon *et al.*, Phys. Rev. Lett. **49**, 978 (1982).

[2] M.R. Anders, S. Shlomo and I. Talmi, in preparation.